

Mathematical tools

M1. Introduction to tensors

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Mathematical tools

M1. Introduction to tensors

1.1 Vector spaces

Definition: set V is called **vector space** over the field $(\mathbb{K}, +, \cdot)$ if:

- the group $(V, +)$ with $+$: $V \times V \rightarrow V$ is abelian
- there exists an operation \cdot : $\mathbb{K} \times V \rightarrow V$ such that:
for all $\lambda, \mu \in \mathbb{K}$ and $v, w \in V$ there is
 - $\lambda \cdot (v + w) = \lambda \cdot v + \lambda \cdot w$,
 - $(\lambda + \mu) \cdot v = \lambda \cdot v + \mu \cdot v$,
 - $\lambda \cdot (\mu \cdot v) = (\lambda \cdot \mu) \cdot v$,
 - $1_{\mathbb{K}} \cdot v = v$.

Example: $\mathbb{K} := \mathbb{R}$, $V := \mathbb{R}^n$ defined by

$$V := \left\{ (v_1, \dots, v_n)^T \mid v_1, \dots, v_n \in \mathbb{R} \right\}$$

$n = \dim(V)$ is the **dimension** of V .

Basis: there exist a **basis** $\mathcal{B} = \langle v_1, v_2, \dots, v_n \rangle$, composed of n linearly independent **basis vectors** $v_i \in V$.

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M1. Introduction to tensors

1.1 Vector spaces: Dual space

Definition: a **scalar product** on V is the application

$(\cdot, \cdot) : V \times V \rightarrow \mathbb{K}$ which, for all $u, v \in V$

- (u, v) is bilinear
- (u, v) is conjugate-symmetric
- (u, u) is positive definite, ie. $(u, u) \geq 0$ and $(u, u) = 0$ implies $u = 0$.

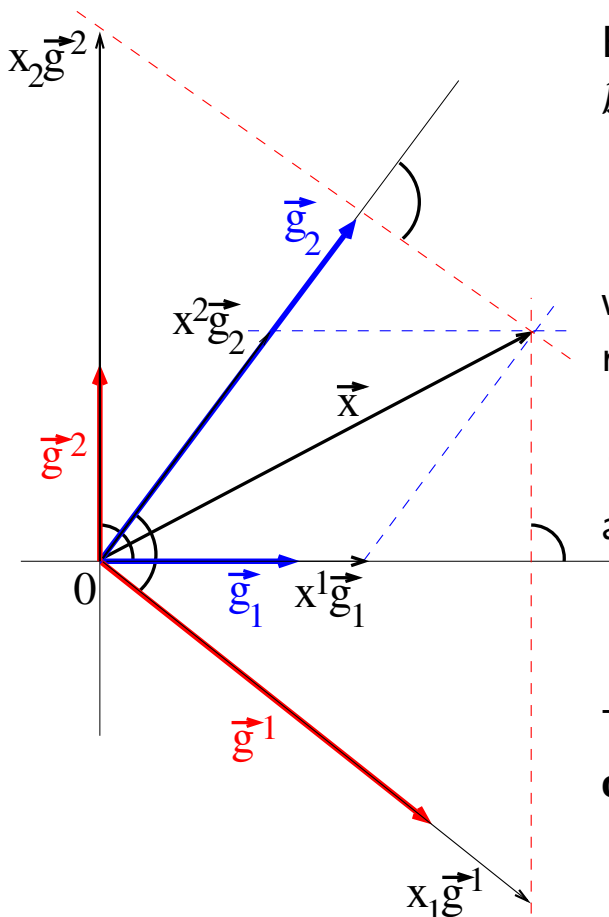
Definition: **euclidean vector-space** is a vector space V together with a scalar product.

Definition: the **dual space** V^* to a vector space V is the set of all linear functions $f : V \rightarrow \mathbb{R}$. For a euclidean space $V^* \cong V$ because of the representation of any $f(u) \in V^*$ by $f(u) = (v, u)$.

Definition: For any vector space V and its basis $\mathcal{B} = \langle v_1, \dots, v_n \rangle$ we can construct a **dual basis** $\mathcal{B}^* = \langle v_1^*, \dots, v_n^* \rangle$ of V^* such that $v_i^*(v_j) = \delta_{ij}$.



1.2 Basis and dual basis for \mathbb{R}^n



Every $\mathbf{x} \in V$ can be represented in the basis $\mathcal{B} = \langle \mathbf{g}_1, \mathbf{g}_2, \dots, \mathbf{g}_n \rangle$ by:

$$\mathbf{x} = x^i \mathbf{g}_i$$

where x^i are **coordinates** of vector \mathbf{x} with respect to the basis $\mathcal{B} = \langle \mathbf{g}_1, \mathbf{g}_2, \dots, \mathbf{g}_n \rangle$

“Other” coordinates of \mathbf{x} can be obtained by a scalar product:

$$x_j = \mathbf{x} \cdot \mathbf{g}_j \quad (\neq x^j \text{ in general!})$$

These coordinates are with respect to the **reciprocal / dual basis** $\mathcal{B}' = \langle \mathbf{g}^1, \mathbf{g}^2, \dots, \mathbf{g}^n \rangle$,

$$\mathbf{x} = x_j \mathbf{g}^j$$



1.2 Basis and dual basis: mutual relations

Starting from

$$\mathbf{x} = x^i \mathbf{g}_i = x_j \mathbf{g}^j \quad \text{where} \quad x_i := \mathbf{x} \cdot \mathbf{g}_i \quad (1)$$

we multiply (scalar product) by \mathbf{g}_k to obtain:

$$x_k = g_{ki} x^i \quad \text{with} \quad g_{ki} := \mathbf{g}_k \cdot \mathbf{g}_i. \quad (2)$$

The matrix g_{ki} is called the **covariant metric tensor** (cf. later).
Substitute (2) back to (1) to obtain:

$$\mathbf{x} = x^i \mathbf{g}_i = x^i g_{ji} \mathbf{g}^j$$

Hence

$$\mathbf{g}_i = g_{ij} \mathbf{g}^j \quad \text{and} \quad \mathbf{g}^k = g^{ki} \mathbf{g}_i, \quad (3)$$

where g^{ki} is the **contravariant metric tensor** defined by:

$$g^{ki} g_{ij} = \delta_j^k \quad \text{i.e. inverse of } g_{ij}.$$



1.2 Basis and dual basis: mutual relations

From (3) and $g_{ij} := \mathbf{g}_i \cdot \mathbf{g}_j$:

$$\mathbf{g}^k \cdot \mathbf{g}_\ell = g^{ki} \mathbf{g}_i \cdot \mathbf{g}_\ell = g^{ki} g_{i\ell} = \delta_\ell^k \quad (4)$$

Hence the dual basis is orthogonal to the basis and from (3)

$$g^{k\ell} = \mathbf{g}^k \cdot \mathbf{g}^\ell.$$

Link between \mathcal{B}' of V and \mathcal{B}^* of V^* : through the scalar product:

$$v_i^*(\mathbf{x}) = (\mathbf{g}^i, \mathbf{x}) = \mathbf{g}^i \cdot \mathbf{x}$$

thanks to (4).



1.3 Covariant and contravariant coordinates of a vector

The role of sub and super-scripts

Basis change through a bijection $F : \mathbb{R}^n \rightarrow \mathbb{R}^n$ (matrix):

Represent F through a matrix $[F_j^i]$.

Introduce a new basis $\tilde{\mathcal{B}} = \langle \tilde{\mathbf{g}}_1, \dots, \tilde{\mathbf{g}}_n \rangle$ through

$$\tilde{\mathbf{g}}_j := F_j^i \mathbf{g}_i \quad \text{and hence} \quad \mathbf{g}_j = \bar{F}_j^k \tilde{\mathbf{g}}_k.$$

where $[\bar{F}_k^j]$ is the inverse of $[F_j^i]$, ie. $\bar{F}_k^j \cdot F_j^i = \delta_k^i$.

Search for coordinates \tilde{x}^ℓ of \mathbf{x} with respect to \mathbf{g}_ℓ :

$$\mathbf{x} = x^i \mathbf{g}_i = x^i \delta_i^k \mathbf{g}_k = x^i \bar{F}_i^\ell F_\ell^k \mathbf{g}_k = \tilde{x}^\ell \tilde{\mathbf{g}}_\ell$$

Hence

$$\tilde{x}^\ell = x^i \bar{F}_i^\ell \quad \text{and} \quad \tilde{x}_\ell = \mathbf{x} \cdot \tilde{\mathbf{g}}_\ell = \mathbf{x} \cdot [F_\ell^i \mathbf{g}_i] = F_\ell^i (\mathbf{x} \cdot \mathbf{g}_i) = F_\ell^i x_i.$$

And also

$$\tilde{\mathbf{g}}^\ell = \bar{F}_k^\ell \mathbf{g}^k$$



1.3 Covariant and contravariant coordinates of a vector

The role of sub and super-scripts

Thus:

$$\tilde{\mathbf{g}}_j = F_j^i \mathbf{g}_i \quad \text{and} \quad \tilde{\mathbf{g}}^\ell = \bar{F}_k^\ell \mathbf{g}^k$$

$$\tilde{x}_\ell = F_\ell^i x_i \quad \text{and} \quad \tilde{x}^\ell = \bar{F}_k^\ell x^k$$

Observation:

- \mathbf{g}_i and x_i transform to $\tilde{\mathbf{g}}_\ell$ and \tilde{x}_ℓ with F_ℓ^i
 \Rightarrow “co-variant” = “with transformation F ”,
- \mathbf{g}^k and x^k transform to $\tilde{\mathbf{g}}^\ell$ and \tilde{x}^ℓ against F
 \Rightarrow “contra-variant” = “against F ” = “through an inverse mapping”.



2.1 Tensor: possible definitions

Let V be a real vector-space and V^* its dual.

Definition (tensor as a multi-linear map): A tensor is a multi-linear application $T : \underbrace{V \times \dots \times V}_{K \text{ copies}} \times \underbrace{V^* \times \dots \times V^*}_{\bar{K} \text{ copies}} \rightarrow \mathbb{R}$.

T is a K -times covariant and \bar{K} -times contravariant tensor of order $N = K + \bar{K}$.

For a euclidean space $V = \mathbb{R}^n$, we can represent V by the basis $\mathcal{B} = \langle \mathbf{g}_1, \dots, \mathbf{g}_n \rangle$ and V^* by the reciprocal basis $\mathcal{B}' = \langle \mathbf{g}^1, \dots, \mathbf{g}^n \rangle$:

Definition (tensor as a hyper-matrix): the representation hyper-matrix \mathbf{T} in \mathbb{R}^n of the multi-linear application T writes

$$\mathbf{T} = T_{i_1 \dots i_K}^{j_1 \dots j_{\bar{K}}} \mathbf{g}^{i_1} \otimes \dots \otimes \mathbf{g}^{i_K} \otimes \mathbf{g}_{j_1} \otimes \dots \otimes \mathbf{g}_{j_{\bar{K}}}$$

where $\{T_{i_1 \dots i_K}^{j_1 \dots j_{\bar{K}}}\}$ are the components of the tensor with respect to the basis $\mathcal{B} = \langle \mathbf{g}_1, \dots, \mathbf{g}_n \rangle$ (depends on this basis!).

2.1 Tensor: possible definitions

Definition (tensor as a $(K + \bar{K})$ -dimensional table):

An ordered $n \times (K + \bar{K})$ -tuple $T_{i_1 \dots i_K}^{j_1 \dots j_{\bar{K}}}$ is a (K -times covariant and \bar{K} -times contravariant) tensor of order $N = K + \bar{K}$ in \mathbb{R}^n , if for any basis-transformation $\tilde{\mathbf{g}}_j = F_j^i \mathbf{g}_i$ in \mathbb{R}^n and its inverse $\mathbf{g}_j = \bar{F}_j^i \tilde{\mathbf{g}}_i$, it transforms by the following recipe

$$\tilde{T}_{\tilde{i}_1 \dots \tilde{i}_K}^{\tilde{j}_1 \dots \tilde{j}_{\bar{K}}} = F_{\tilde{i}_1}^{i_1} \dots F_{\tilde{i}_K}^{i_K} \cdot \bar{F}_{j_1}^{\tilde{j}_1} \dots \bar{F}_{j_{\bar{K}}}^{\tilde{j}_{\bar{K}}} \cdot T_{i_1 \dots i_K}^{j_1 \dots j_{\bar{K}}}$$

Link between the definitions:

Let us evaluate $T(\mathbf{x}_1, \dots, \mathbf{x}_K, \mathbf{y}_1, \dots, \mathbf{y}_{\bar{K}})$ for a special choice

$$\mathbf{x}_\ell = \mathbf{g}_{i_\ell} \quad \text{and} \quad \mathbf{y}_\ell = \mathbf{g}^{j_\ell}.$$

The multi-linear map gives

$$T(\mathbf{g}_{i_1}, \dots, \mathbf{g}_{i_K}, \mathbf{g}^{j_1}, \dots, \mathbf{g}^{j_{\bar{K}}}) = T_{i_1 \dots i_K}^{j_1 \dots j_{\bar{K}}}.$$

For $\mathbf{x}_\ell = x_{,\ell}^i \mathbf{g}_i$ and $\mathbf{y}_k = y_{j,k} \mathbf{g}^j$ by multi-linearity

$$T(\mathbf{x}_1, \dots, \mathbf{x}_K, \mathbf{y}_1, \dots, \mathbf{y}_{\bar{K}}) = T_{i_1 \dots i_K}^{j_1 \dots j_{\bar{K}}} x_{,1}^{i_1} \dots x_{,K}^{i_K} \cdot y_{j_1,1} \dots y_{j_{\bar{K}},\bar{K}}$$

2.2 Basic operations on tensors (on tensor components)

Let \mathbf{A} , \mathbf{B} be tensors (hyper-matrices).

- **addition** of tensors of the same type:

$$\mathbf{C} = \mathbf{A} + \mathbf{B}, \quad \text{e.g. } C_{ij} = A_{ij} + B_{ij},$$

- **multiplication by** $\alpha \in \mathbb{R}$:

$$\mathbf{C} = \alpha \mathbf{A}, \quad \text{e.g. } C_{ij} = \alpha A_{ij},$$

- **tensor-multiplication**:

$$\mathbf{C} = \mathbf{A} \otimes \mathbf{B}, \quad \text{e.g. } C_{kl}^{ij} = A^{ij} \cdot B_{kl},$$

- **contraction of indices**:

$$\mathbf{C} = \mathbf{A}\mathbf{B}, \quad \text{e.g. } C_k^\ell = A_{ki} B^{i\ell},$$

$$\gamma = \mathbf{A} : \mathbf{B}, \quad \text{e.g. } \gamma = A_{ij} B^{ij},$$

- **raising or lowering of an index**: contravariant \leftrightarrow covariant component:

$$C_k^i = C_{jk} g^{ij} \quad \text{or} \quad C_{ikl} = C_{kl}^j g_{ij},$$



2.3 Special operations for 2nd-order tensors

Let \mathbf{A} be a 2nd-order tensor (hyper-matrix)

$$\mathbf{A} = A^{ij} \mathbf{g}_i \otimes \mathbf{g}_j = A_{ij} \mathbf{g}^i \otimes \mathbf{g}^j = A_j^i \mathbf{g}_i \otimes \mathbf{g}^j = A_i^j \mathbf{g}^i \otimes \mathbf{g}_j$$

- **transpose**: $\mathbf{B} = \mathbf{A}^T \quad B_{ji} = A_{ij}, \quad B_{\cdot i}^j = A_{i \cdot}^j.$

- **trace**: $\text{tr}(\mathbf{A}) = A_i^i = g^{ij} A_{ij} = g_{ij} A^{ij}$

- **Frobenius norm**: $|\mathbf{A}| = \sqrt{\mathbf{A} : \mathbf{A}} = \text{tr}(\mathbf{A}^T \mathbf{A}) = \text{tr}(\mathbf{A} \mathbf{A}^T) = (A_{ij} A^{ij})^{1/2}$

- **determinant**: $\det(\mathbf{A})$

$$\det(\mathbf{A}) = \det[A_j^i] = \det[A_i^j] = \det[g_{ij}] \det[A^{ij}] = \det[g^{ij}] \det[A_{ij}].$$



2.3 Decompositions of a 2nd-order tensor

- **symmetric and skew (antisymmetric) tensors:** $\mathbf{A} = \mathbf{S} + \mathbf{W}$,

$$\mathbf{S} = \frac{1}{2} (\mathbf{A} + \mathbf{A}^T) \quad \text{and} \quad \mathbf{W} = \frac{1}{2} (\mathbf{A} - \mathbf{A}^T).$$

- $\mathbf{S} = \mathbf{S}^T$, $\mathbf{W} = -\mathbf{W}^T$.

- For every \mathbf{B} (2nd order tensor):

$$\mathbf{S} : \mathbf{B} = \mathbf{S} : \mathbf{B}^T = \mathbf{S} : \frac{1}{2} (\mathbf{B} + \mathbf{B}^T)$$

$$\mathbf{W} : \mathbf{B} = -\mathbf{W} : \mathbf{B}^T = \mathbf{W} : \frac{1}{2} (\mathbf{B} - \mathbf{B}^T)$$

$$\mathbf{S} : \mathbf{W} = 0$$

- $\exists \mathbf{w} \in \mathbb{R}^n$ s.t. $\forall \mathbf{u} \in \mathbb{R}^n$ there is $\mathbf{W}\mathbf{u} = \mathbf{w} \times \mathbf{u}$.

2.3 Decompositions of a 2nd-order tensor

- **spherical part $\alpha \mathbf{I}$ and deviator $\text{dev}(\mathbf{A})$ of \mathbf{A} ,**

$$\mathbf{A} = \alpha \mathbf{I} + \text{dev}(\mathbf{A}) \quad \text{with} \quad \alpha := \frac{1}{n} \text{tr}(\mathbf{A}) = \frac{1}{n} (\mathbf{I} : \mathbf{A})$$

- $\text{tr}(\text{dev}(\mathbf{A})) = 0$

2.3 Eigenvalues, eigenvectors and principal scalar invariants

Eigenvalue problem: find all pairs $(\mathbf{u}, \lambda) \in \mathbb{C}^n \times \mathbb{C}$ such that

$$\mathbf{A}\mathbf{u} = \lambda\mathbf{u}. \quad (5)$$

Let \mathbf{Q} is an orthogonal matrix (2nd-order tensor), then from (5)

$$\left(\mathbf{Q}^T \mathbf{A} \mathbf{Q}\right) \mathbf{Q}^T \mathbf{u} = \lambda \mathbf{Q}^T \mathbf{u},$$

ie. eigenvalues of \mathbf{A} are equal to eigenvalues of $\mathbf{Q}^T \mathbf{A} \mathbf{Q}$.



2.3 Eigenvalues, eigenvectors and principal scalar invariants

Eigenvalues of \mathbf{A} are invariant to rotations \Rightarrow characteristic polynomial is also invariant

$$\begin{aligned} \det[\mathbf{A} - \lambda \mathbf{I}] &= 0 \\ \lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 &= 0 \end{aligned}$$

with the following principal scalar invariants:

$$I_1(\mathbf{A}) = \text{tr}(\mathbf{A})$$

$$I_2(\mathbf{A}) = \frac{1}{2} [\text{tr}^2(\mathbf{A}) - \text{tr}(\mathbf{A}^2)] = \text{tr}(\mathbf{A}^{-1}) \det(\mathbf{A})$$

$$I_3(\mathbf{A}) = \det(\mathbf{A})$$



Cayley-Hamilton theorem: for every 2nd-order tensor \mathbf{A} with its characteristic equation

$$\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0,$$

there is

$$\mathbf{A}^3 - I_1 \mathbf{A}^2 + I_2 \mathbf{A} - I_3 \mathbf{I} = \mathbf{0},$$

3.1 Levi-Civita permutation symbol

Definition: Levi-Civita permutation symbol e_{ijk} and e^{ijk} :

$$e_{ijk} = \begin{cases} 1 & \text{if } (i, j, k) \text{ is an even permutation of } (1, 2, 3), \\ -1 & \text{if } (i, j, k) \text{ is an odd permutation of } (1, 2, 3), \\ 0 & \text{otherwise } (i = j \text{ or } i = k \text{ or } j = k). \end{cases}$$

$$e^{ijk} = e_{ijk} \quad , \quad e^{ijk} e_{ijk} = 6.$$

Matrix determinant from Linear Algebra I+II:

Determinant of a matrix $A \in \mathbb{R}_3^{3 \times 3}$ was defined by:

$$\det(A) = \sum_{\sigma \in S_3} \text{sign}(\sigma) \prod_{i=1}^3 A_{i, \sigma(i)} \quad , \quad A = [A_{ij}].$$

Here, S_3 is the set of all permutations of $\{1, 2, 3\}$ and $\text{sign}(\sigma) = +1$ for even permutations σ and -1 for odd ones.

Determinant with Levi-Civita permutation symbol:

$$\det(A) = e^{ijk} A_{1i} A_{2j} A_{3k} = \frac{1}{6} e^{ijk} e^{\ell mn} A_{\ell i} A_{mj} A_{nk}.$$

But attention, Levi-Civita permutation symbol is not a tensor!

3.2 Levi-Civita permutation tensor

Levi-Civita permutation symbol is not a tensor:

proof by change of basis $\tilde{\mathbf{g}}_j = F_j^i \mathbf{g}_i$:

$$F_i^\ell F_j^m F_k^n e_{\ell mn} = \det(F) e_{ijk} \neq e_{ijk}!$$

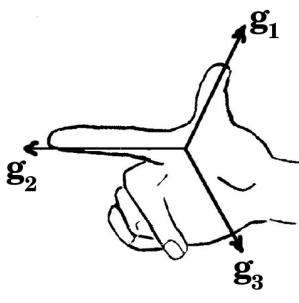
However, Levi-Civita symbol can form the following tensor.

Definition: Levi-Civita permutation tensor ε_{ijk} and ε^{ijk} :

$$\varepsilon_{ijk} = \Delta_g \sqrt{g} e_{ijk} \quad , \quad \varepsilon^{ijk} = \frac{\Delta_g}{\sqrt{g}} e^{ijk}$$

$$g = \det[g_{ij}] \quad , \quad \Delta_g = \begin{cases} +1 & \text{if } \langle \mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3 \rangle \text{ right-handed,} \\ -1 & \text{if } \langle \mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3 \rangle \text{ left-handed.} \end{cases}$$

Right-handedness of a basis (coordinate system) $\langle \mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3 \rangle$:



right-handed



???



3.2 Levi-Civita permutation tensor

Proof of tensor-like behaviour: change of basis $\tilde{\mathbf{g}}_j = F_j^i \mathbf{g}_i$:

$$\begin{aligned} \tilde{\varepsilon}_{ijk} &= F_i^\ell F_j^m F_k^n \varepsilon_{\ell mn} = F_i^\ell F_j^m F_k^n \Delta_g \sqrt{g} e_{\ell mn} \\ &= \Delta_g \sqrt{\det[g_{ij}]} \det(F) e_{ijk} \\ &= \sqrt{\frac{\det[\tilde{g}_{ij}]}{\det^2(F)}} \Delta_g \det(F) e_{ijk} \\ &= \sqrt{\det[\tilde{g}_{ij}]} \Delta_g \text{sign}(\det(F)) e_{ijk} \\ &= \sqrt{\tilde{g}} \Delta_{\tilde{g}} e_{ijk} \end{aligned}$$

Hence, the tensor-like transformation of covariant components $\varepsilon_{\ell mn}$ before basis-change (above) gives exactly the covariant components $\tilde{\varepsilon}_{ijk}$ of Levi-Civita tensor with respect to the transformed basis (below).

3.3 Levi-Civita permutation tensor: volume of triad $\mathbf{a}, \mathbf{b}, \mathbf{c}$

Let $V : \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{R}^3 \rightarrow \mathbb{R}$ be defined for all $\mathbf{a}, \mathbf{b}, \mathbf{c} \in \mathbb{R}^3$ by

$$V(\mathbf{a}, \mathbf{b}, \mathbf{c}) = (\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}$$

This application is multi-linear \rightarrow function V defines a tensor.

Let us express \mathbf{a}, \mathbf{b} and \mathbf{c} with respect to a basis $\langle \mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3 \rangle$:

$$\mathbf{a} = a^i \mathbf{g}_i \quad , \quad \mathbf{b} = b^j \mathbf{g}_j \quad , \quad \mathbf{c} = c^k \mathbf{g}_k$$

The $3 \times$ covariant components ϑ_{ijk} of $\underline{\underline{V}}$ must satisfy

$$V(\mathbf{a}, \mathbf{b}, \mathbf{c}) = \vartheta_{ijk} a^i b^j c^k$$

It appears that $\vartheta_{ijk} \equiv \varepsilon_{ijk}$, ie. $\underline{\underline{V}}$ coincides with the Levi-Civita permutation tensor! Prove it! Idea: verify that it is so for canonical coordinates and use unicity of tensors for other bases



3.3 Levi-Civita permutation tensor: vector product

Let us denote $\mathbf{d} = \mathbf{a} \times \mathbf{b}$. Then

$$V(\mathbf{a}, \mathbf{b}, \mathbf{c}) = (\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c} = \mathbf{d} \cdot \mathbf{c}$$

We know that scalar product is also multi-linear, hence tensor

$$\mathbf{d} \cdot \mathbf{c} = d_k c^k = d^j c_j.$$

We have just shown that

$$V(\mathbf{a}, \mathbf{b}, \mathbf{c}) = \varepsilon_{ijk} a^i b^j c^k = d_k c^k$$

Hence, covariant components of vector product can be easily computed at any coordinate system by

$$d_k = \varepsilon_{ijk} a^i b^j \quad \text{and then} \quad d^i = g^{ij} d_j$$



Cramer's rule: Levi-Civita tensor vs. metric tensor

$$g^{kn} = \frac{1}{2} \varepsilon^{ijk} \varepsilon^{\ell mn} g_{il} g_{jm} \quad , \quad g_{kn} = \frac{1}{2} \varepsilon_{ijk} \varepsilon_{\ell mn} g^{il} g^{jm}$$

Cramer's rule to solve linear systems $[g^{ij}] = [g_{kl}]^{-1}$!